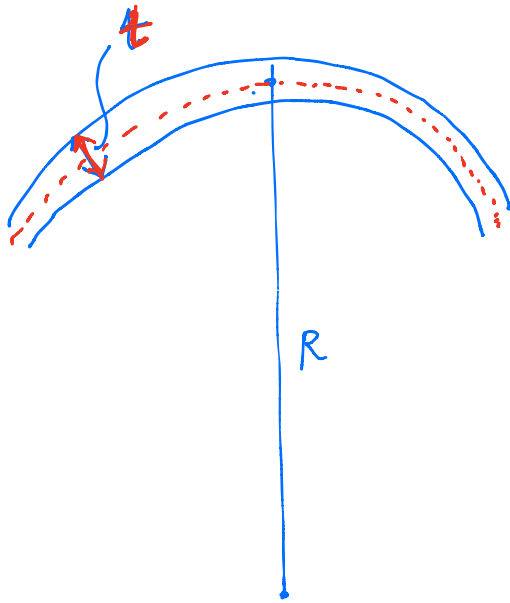
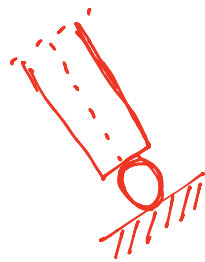


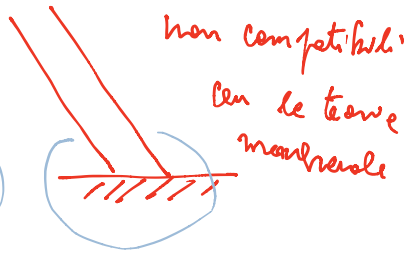
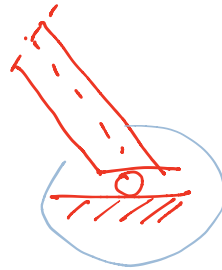
GUSCI - TEORIA DEI GUSCI SOTTILI



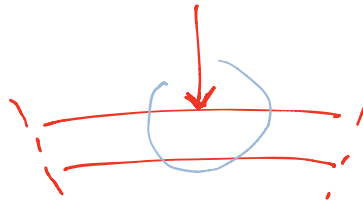
$$\frac{t}{R} \ll \frac{1}{50} \Rightarrow \text{guscio sottile}$$



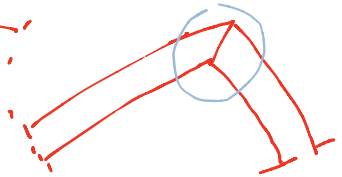
Compatibile con le
teorie membranale



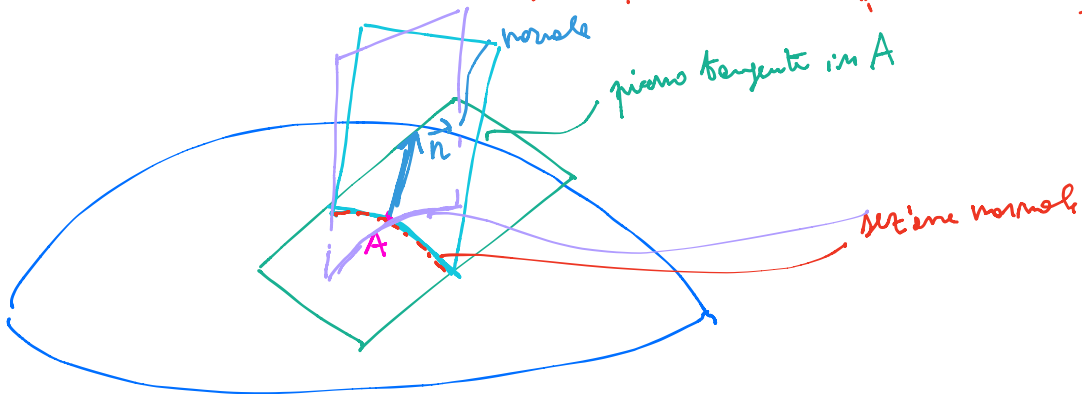
non compatibile
con le teorie
membranale



non è compatibile

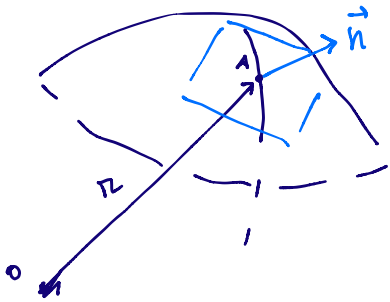


CLASSIFICAZIONE DEI GUSCI



A l'ordine primo normale corrisponde una curvatura locale k e un raggio di curvatura r

$$k = \pm \frac{1}{r}$$



Se l'origine del raggio di curvatura è positiva rispetto a \vec{n}

allora $k = \frac{1}{r}$

o

negative rispetto a \vec{n}

Ci saranno (ovviamente) un

allora $k = -\frac{1}{r}$

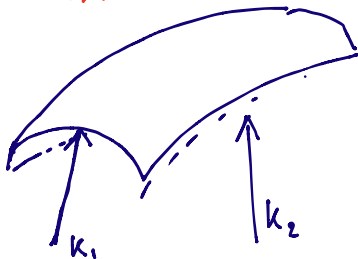
valore minimo e un valore massimo di k in A

k_1 e k_2 curvatura principali

Le due direzioni principali (corrispondenti a k_1 e k_2) sono ortogonali

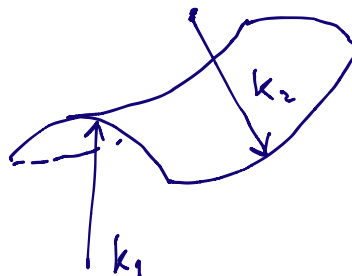
CURVATURA GAUSSIANA $k_g = k_1 \cdot k_2$

similettico



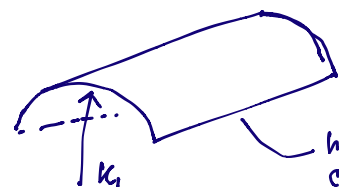
$$k_g = k_1 \cdot k_2 > 0$$

anti-elastico



$$k_g = k_1 \cdot k_2 < 0$$

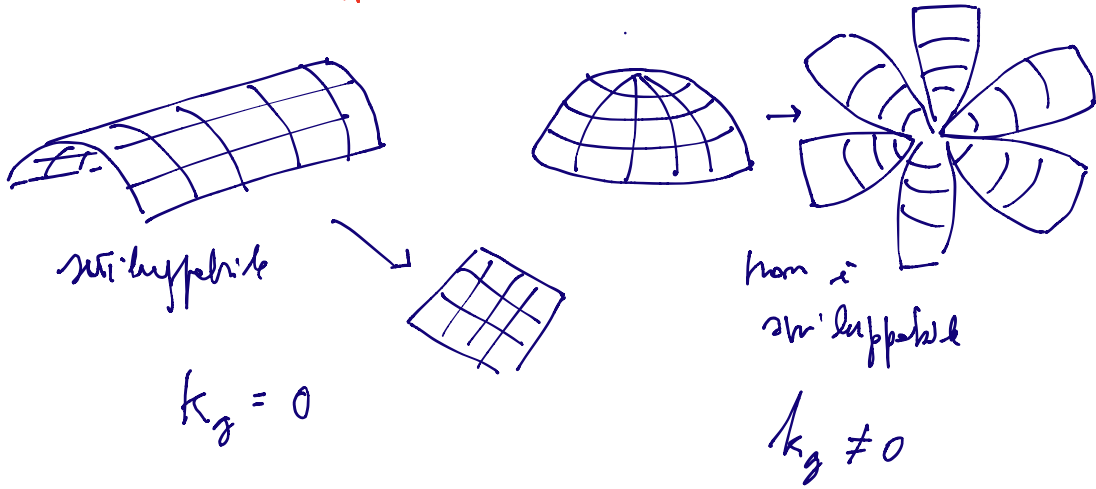
a singola curvatura



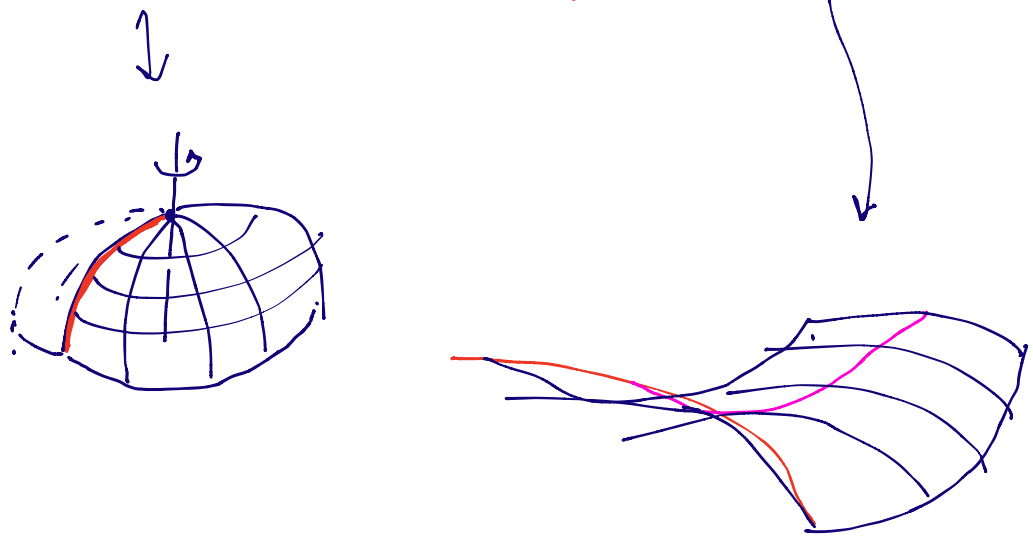
$$k_g = 0$$

non c'è curvatura

Superfici sviluppabili e non-sviluppabili

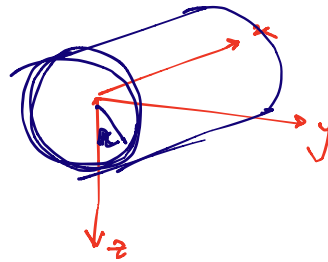


Superfici di rivoluzione, superficie di trasferimento



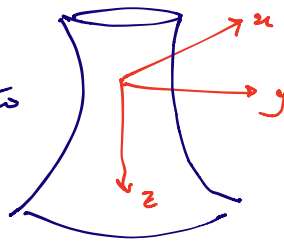
DESCRIZIONE ANALITICA DEI GUSCI (DELLA SUPERFICIE MEDIA)

$$x^2 + y^2 = r^2 \quad \text{cilindro}$$

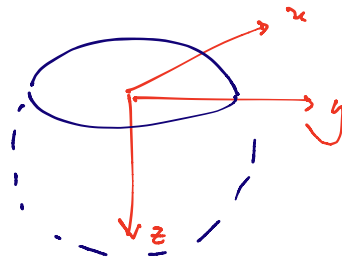


$$\frac{x^2 + y^2}{a^2} - \frac{z^2}{b^2} = 1$$

toro
di
raffreddamento

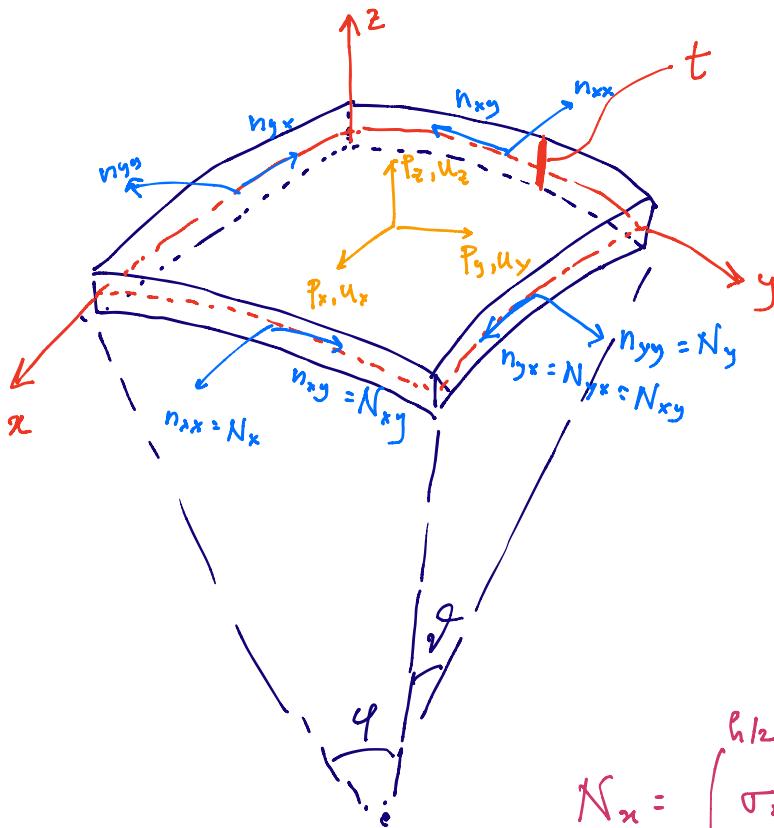


$$\frac{x^2 + y^2 + z^2}{a^2} = 1$$



TEORIA MEMBRANALE PER GUSCI CON CURVATURE

PRINCIPALI



$$k_1 < 0 \text{ e } k_2 < 0$$

$$\vec{u} = (u_x, u_y, u_z)^t$$

$$\vec{p} = (p_x, p_y, p_z)^t$$

$$\vec{\epsilon} = (\epsilon_x, \epsilon_y, \gamma_{xy})^t$$

$$\vec{\sigma} = (\sigma_x, \sigma_y, \tau_{xy})^t$$

$$\vec{S} = (N_x, N_y, N_{xy})^t$$

sul libro

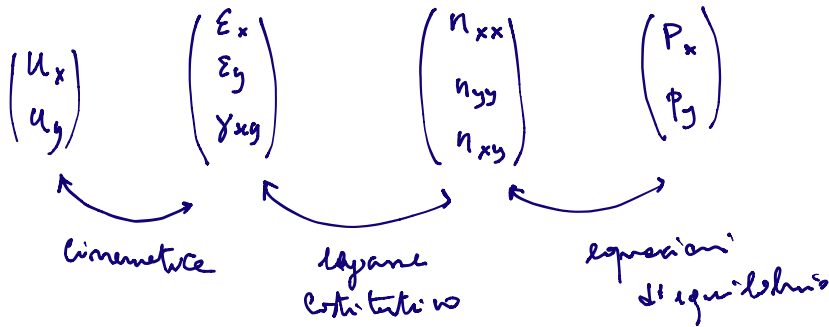
$$n_{xx} \quad n_{yy} \quad n_{xy}$$

$$N_x = \int_{-h/2}^{h/2} \sigma_x dz \quad \text{fun. de. piastra}$$

ADesso DOVREMO TENERE CONTO DELLA CURVATURA

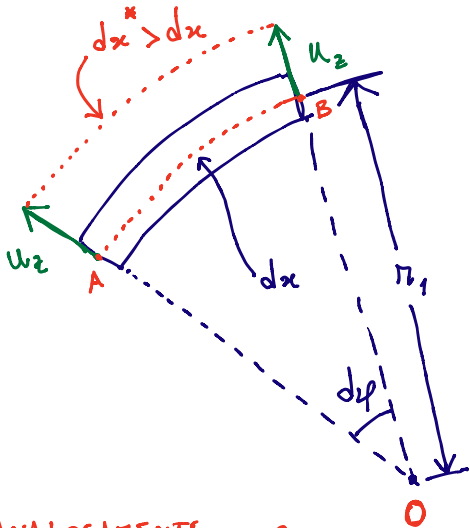
$$N_z = \int_{-t/2}^{t/2} \sigma_x \left(1 - \frac{z}{r_2}\right) dz \approx \int_{-t/2}^{t/2} \sigma_x dz \quad \begin{array}{l} \text{TEORIA} \\ \text{MEMBRANALE} \end{array}$$

per gusci sottili e trascurabile



CINEMATICS

$$\left[\epsilon_x = \frac{\partial u_x}{\partial x} ; \epsilon_y = \frac{\partial u_y}{\partial y} ; \gamma_{xy} = \frac{\partial u_x}{\partial y} + \frac{\partial u_y}{\partial x} \right] \text{ per una piastra}$$



$$dx = r_1 \sin(d\varphi) \approx r_1 d\varphi$$

$$\epsilon_x = \frac{dx^* - dx}{dx} = \frac{(r_1 + u_z) d\varphi - r_1 d\varphi}{r_1 d\varphi}$$

$$= \frac{u_z}{r_1} = -k_1 u_z = \epsilon_x$$

ANALOGAMENTE per ϵ_y

$$\epsilon_y = \frac{u_z}{r_2} = -k_2 u_z$$

$$\begin{pmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \partial/\partial x & 0 & -k_1 \\ 0 & \partial/\partial y & -k_2 \\ \partial/\partial y & \partial/\partial x & 0 \end{pmatrix} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix}, \quad \vec{\varepsilon} = \underline{B} \vec{u}$$

operatore di flessione

LEGAME COSTITUTIVO

Considerando che il materiale del ^{quello} pie è elastico lineare e quindi, si è governato da una legge alla Hooke

$$\begin{pmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{pmatrix} = \frac{E}{1-\nu^2} \begin{pmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{pmatrix} \begin{pmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{pmatrix}$$

$$\begin{pmatrix} n_{xx} \\ n_{yy} \\ n_{xy} \end{pmatrix} = \frac{E t}{1-\nu^2} \begin{pmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{pmatrix} \begin{pmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{pmatrix}$$

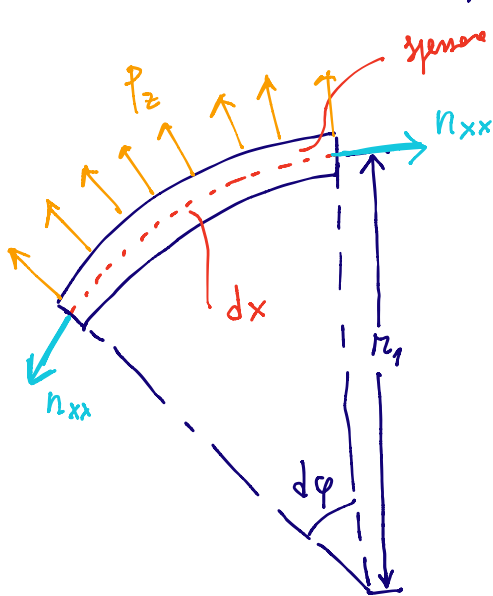
$$n_{ij} = \int_{-t/2}^{t/2} \sigma_{ij} \left(1 - \frac{z}{\pi}\right) dz \approx \int_{-t/2}^{t/2} \sigma_{ij} dz = \begin{matrix} \text{per questi} \\ \text{materiali} \end{matrix} \begin{matrix} \text{isotropici} \\ \text{e omogenei} \\ \text{costante} \end{matrix} = t \sigma_{ij}$$

$$\underline{D} = \frac{E t}{1-\nu^2} \begin{pmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{pmatrix} \begin{matrix} \text{MATRICE} \\ \text{DI} \\ \text{MOIDENZA} \end{matrix} \quad \vec{S} = \underline{D} \vec{\varepsilon}$$

EQUAZIONI D'EQUILIBRIO

$$\begin{cases} \frac{\partial n_{xx}}{\partial x} + \frac{\partial n_{xy}}{\partial y} + \phi_x = 0 \\ \frac{\partial n_{xy}}{\partial x} + \frac{\partial n_{yy}}{\partial y} + \phi_y = 0 \end{cases}$$

+ un'altra espressione di tangente conto di n_x e p_z

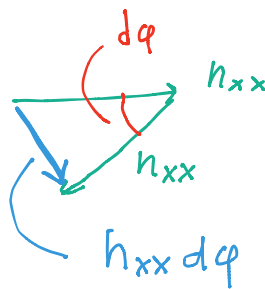


spessore unitario

\Downarrow
 k_1 e k_2

$$dx = r_1 d\varphi$$

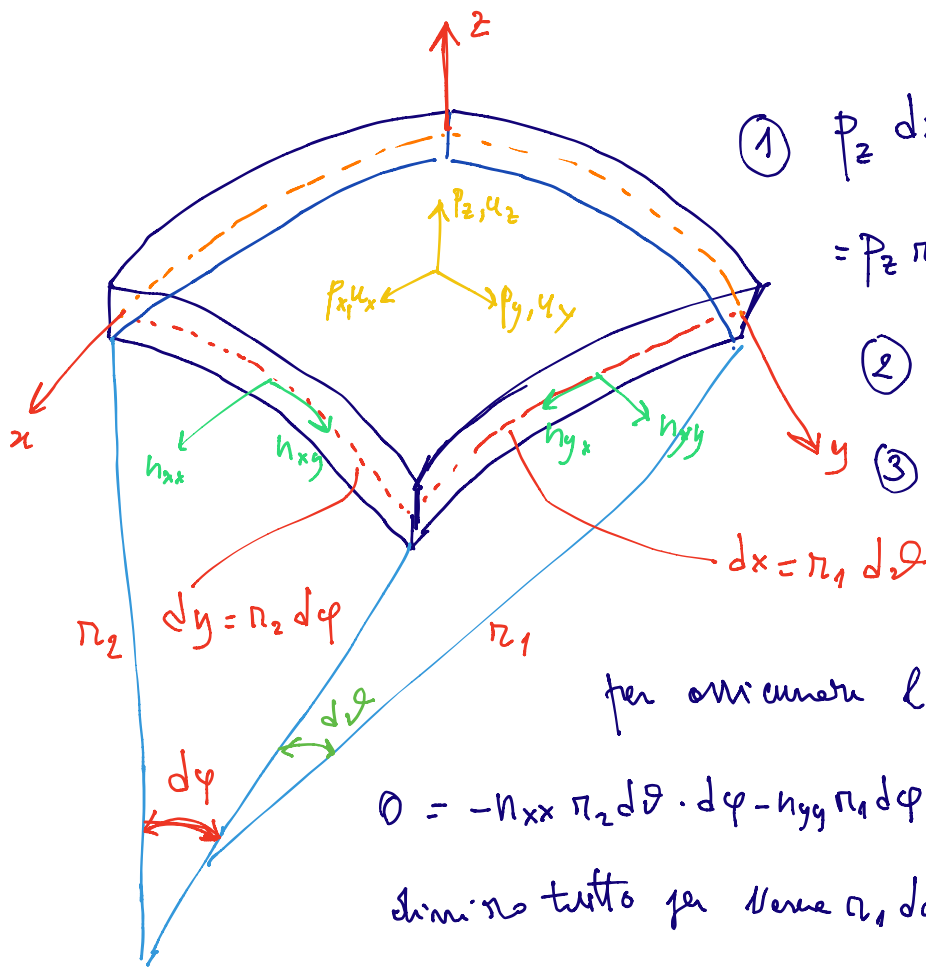
$$p_z dx = p_z r_1 d\varphi$$



$-n_{xx} d\varphi + p_z r_1 d\varphi = 0$
dividiamo tutto per $dx = r_1 d\varphi$

$$-\frac{n_{xx}}{r_1} + p_z = 0 \Rightarrow k_1 n_{xx} + p_z = 0$$

$$k_2 n_{yy} + p_z = 0$$



- ① $p_z dx dy =$
- $= p_z r_1 d\varphi r_2 d\varphi$
- ② $-n_{xx} r_1 d\varphi$
- ③ $-n_{yy} r_2 d\varphi$

per ottenere l'equazione

$$0 = -n_{xx} r_2 d\varphi \cdot d\varphi - n_{yy} r_1 d\varphi \cdot d\varphi + p_z r_1 d\varphi r_2 d\varphi$$

dividiamo tutto per $r_1 d\varphi \cdot r_2 d\varphi$

$$-\frac{n_{xx}}{r_1} - \frac{n_{yy}}{r_2} + p_z = 0$$

$$k_1 n_{xx} + k_2 n_{yy} + p_z = 0$$

$$\left\{ \begin{array}{l} \frac{\partial n_{xx}}{\partial x} + \frac{\partial n_{xy}}{\partial y} + p_x = 0 \\ \frac{\partial n_{xy}}{\partial x} + \frac{\partial n_{yy}}{\partial y} + p_y = 0 \\ k_1 n_{xx} + k_2 n_{yy} + p_z = 0 \end{array} \right.$$

in forme matriciale

$$\underline{B}^* \vec{S} = \vec{P}$$

$$\underline{B}^* = \begin{pmatrix} -\frac{\partial}{\partial x} & 0 & -\frac{\partial}{\partial y} \\ 0 & -\frac{\partial}{\partial y} & -\frac{\partial}{\partial x} \\ -k_1 & -k_2 & 0 \end{pmatrix}$$